





Ultrafast light-driven simultaneous excitation of coherent terahertz magnons and phonons in multiferroic BiFeO₃

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The ultrafast switching of magnetization in multiferroic materials by a femtosecond laser could provide various advantages in photonics and magnonics. An efficient approach to control the light-matter interaction is the modulation of ultrafast coherent magnons and phonons in the high-frequency range. Spontaneous Raman and infrared spectra reveal the excitation of magnons and optical phonons in multiferroic BiFeO₃ in the sub-few-terahertz range. However, coherent control of such quasiparticles has not been achieved yet. In this study, we demonstrate that linearly polarized laser pulses simultaneously excite coherent magnons [out-of-plane (Ψ) and in-plane (Φ) cyclone modes] and optical phonon (E mode) in BiFeO₃. Experimental results in conjugation with phenomenological theory, by considering three uniformly distributed magnetic domains, reveal that impulsive stimulated Raman scattering is responsible for the generation of coherent magnons and phonons in BiFeO₃. The observation of these terahertz magnon and optical phonon modes paves the way for the development of ultrafast magnetoelectro-optical devices.

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I. INTRODUCTION

The multiferroics, unique materials simultaneously exhibiting ferroelectric and (anti)ferromagnetic orders [1–3], are promising candidates for the fabrication of next-generation spintronic devices with added functionalities [4]. Recently, oxide multiferroics have attracted considerable attention owing to the magnetoelectric coupling between the orders by manipulating magnetic properties by electric field [5] and ferroelectric polarization by magnetic field [6]. BiFeO₃ stands out as a canonical material in the multiferroic magnetoelectric research because it exhibits ferroelectricity and antiferromagnetism at room temperature [7–11]. Spontaneous Raman [12–17] and infrared (IR) [18–26] spectroscopic measurements on BiFeO₃ have revealed the excitation of magnons and phonons. BiFeO₃ also exhibits nontrivial light-matter interaction, which yields exotic phenomena, such as visible-light photovoltaic effect [27,28], photoinduced ultrafast magnetic-field-dependent terahertz (THz) emission [29], and giant electrochromic behavior [30]. Recently, ultrafast pump-probe spectroscopy has been successfully used to investigate the excitation of coherent magnons and acoustic phonons in BiFeO₃ in the range below 100 GHz [31,32]. Although such observations are quite enticing, excitations of coherent magnon and optical phonon modes in the sub-few-THz regime have not been reported. For practical application, the development of ultrafast magnetoelectro-optical devices requires the generation of optically controlled polarization-sensitive magnon and

phonon modes at considerably higher frequencies of the THz regime.

In this study, femtosecond pump-probe spectroscopy was employed to generate high-frequency ultrafast coherent magnons and optical phonon in a BiFeO₃ single crystal and to understand their excitation and detection mechanisms. We obtained magnon modes in sub-THz frequency range whereas a phonon mode was observed in few-THz range. The signed amplitude of magnons and phonon exhibit sinusoidal dependence on the pump azimuth angle, which is originated from impulsive stimulated Raman scattering (ISRS). Temperature-dependent pump-probe measurements further indicate that the magnon modes of 530, 560, and 740 GHz (at 300 K) belong to $\Psi_1^{(2)}$, $\Psi_1^{(1)}$, and $\Phi_2^{(1,2)}$ symmetries, respectively.

II. MATERIAL AND METHODS

A. Sample fabrication and characteristics

The experiments were performed on a 150- μm -thick single-crystal BiFeO₃, grown by using a modified floating-zone method with laser diodes, parallel to the $\hat{z} \parallel [111]_{\text{pc}}$ orientation [33]. At room temperature, BiFeO₃ exhibits a rhombohedral $R3c$ distorted perovskite structure with a magnetic point group of $3m$. In addition, it exhibits a G -type antiferromagnetic order below the Néel temperature of 640 K [34] with a cycloidal spiral structure having a large modulation period of 62 nm [35,36]. BiFeO₃ is high-temperature ferroelectric with a Curie temperature of 1100 K and spontaneous electric polarization of approximately 100 $\mu\text{C}/\text{cm}^2$ along the z direction [37]. We chose a ferroelectric single domain where the electric polarization is oriented along the sample normal (z axis). The cycloidal rotation was assumed to be along the

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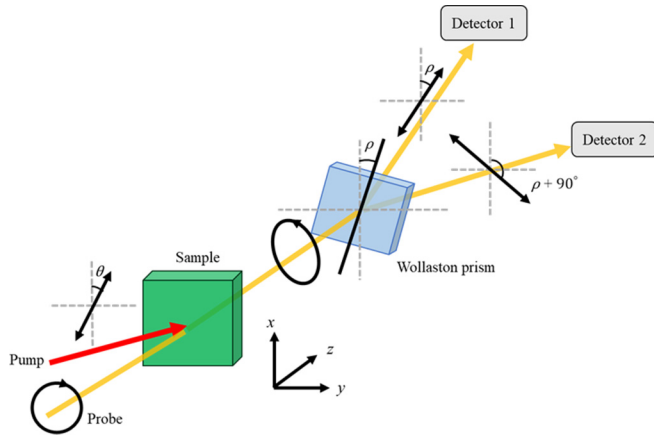


FIG. 1. Schematic of sample orientation pump-probe measurement with a linearly polarized pump and circularly polarized probe. Here θ and ρ are the pump azimuth and the detection angle, respectively.

$\hat{x} \parallel [10\bar{1}]_{\text{pc}}$ direction, whereas $\hat{y} \parallel [\bar{1}2\bar{1}]_{\text{pc}}$ was perpendicular to the plane of cycloidal rotation.

B. Ultrafast pump-probe measurements

To observe the ultrafast spin and lattice dynamics in BiFeO_3 , we performed femtosecond time-resolved pump-probe measurements in the transmission geometry. A pump pulse generated by a Ti:sapphire laser with a regenerative amplifier and optical parametric amplifier (Spectra-Physics) at a repetition rate of 500 Hz, wavelength of 1300 nm, and duration of 70 fs was incident at an angle of 10° with respect to the sample normal (z axis). The pump photon energy is below the lowest $\text{Fe}^{3+}d-d$ or $p-d$ transition energy [38–40]. Probe pulses with a wavelength of 800 nm and repetition rate of 1 kHz were incident perpendicularly to the sample. The linearly polarized pump beam was focused to a spot diameter of approximately $100 \mu\text{m}$, which yielded a fluence of approximately $25 \text{ mJ}/\text{cm}^2$. The circularly polarized probe beam had approximately twice smaller spot diameter. A schematic of the pump-probe measurement and sample orientation are

shown in Fig. 1. In the pump-probe measurement, optically excited coherent magnons and phonons were observed by using the polarimetric detection technique [41–43] with balanced detectors enabling the extraction of the contributions of an induced change in ellipticity of the probe polarization. The pump azimuth angle θ was defined as the angle between the x axis and the direction of the pump polarization (electric field). The detection was performed with a circularly polarized probe in two different geometric configurations, X- and Y-probe configurations [42], where the detection angles were set to $\rho = 0$ and $\pi/4$, respectively.

III. RESULTS AND DISCUSSION

Figure 2 shows the time evolution of the ellipticity change ($\Delta\eta$) in probe polarization in BiFeO_3 at 300 K, for an excitation at a pump azimuth angle of $\theta = -\pi/4$ and detection in the Y-probe configuration. We observed oscillatory signals having different characteristic frequencies, as shown in Figs. 2(a) and 2(b), with an abruptness at delay $t \sim 2.8$ ps. The abrupt signal is observed because of the overlap of the probe pulse with the pump pulse, which is reflected from the back surface of the sample, counterpropagated, and reflected again from the front surface [41]. We performed fast Fourier transform (FFT) of the oscillatory signals at positive delays. The resulting amplitude spectra are shown in the insets of Figs. 2(a) and 2(b). The FFT spectra reveal a high-amplitude mode at 2.4 THz in Fig. 2(a) and three low-amplitude modes at 530, 560, and 740 GHz in Fig. 2(b).

Experimental results in conjugation with the symmetry consideration reveal that the ISRS is responsible for the excitation of the 2.4-THz phonon mode having E symmetry, which is consistent with the mode in previous reports observed by spontaneous Raman [16] and IR [18] spectroscopies. Detailed results and analysis on the selection rule for the excitation and detection of the coherent phonon are presented in Figs. S1 and S2 of Supplemental Material (SM) [44] (see also Refs. [16,45] therein).

The temperature dependence of the 2.4-THz frequency is shown in Fig. 3. The FFT spectra of the oscillation curves corresponding to the 2.4-THz mode are shown in Fig. 3(a).

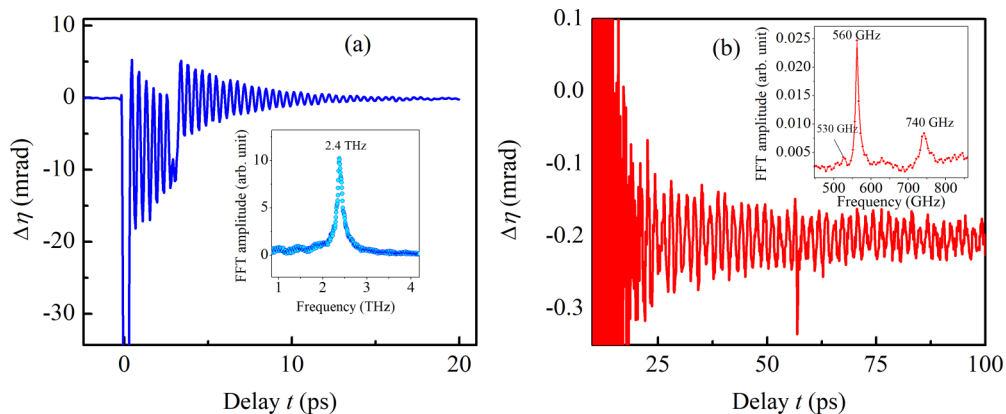


FIG. 2. Temporal evolutions of the ellipticity changes ($\Delta\eta$) in transmitted probe polarizations in BiFeO_3 ($T = 300 \text{ K}$) at two different probe delays up to (a) 20 and (b) 60 ps. The insets show the FFT amplitude spectra of the signals at the corresponding delays, which reveal center frequencies of (a) 2.4 THz and (b) 530, 560, and 740 GHz.

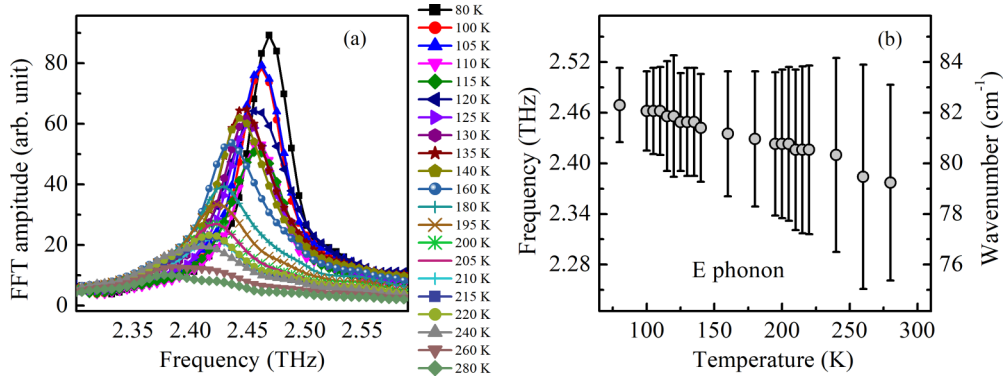


FIG. 3. (a) FFT amplitude spectra of the 2.4-THz mode at various temperatures. (b) Variation of E phonon (2.4 THz) at different temperatures. Here, the black line represents the FWHM of the phonon mode.

The ellipticity curves associated with THz mode is shown in Fig. S3 of SM [44]. Figure 3(b) reveals that the 2.4-THz phonon mode exhibits frequency softening ($\sim 4\%$) within the full width at half maximum (FWHM) of the FFT spectra shown in Fig. 3(a), when the temperature is increased.

To understand the temperature dependence of sub-THz modes, ellipticity changes are recorded at various temperatures as shown in Fig. S4 of SM [44]. The corresponding FFT spectra of the sub-THz frequencies at different temperatures are shown in Fig. 4(a). The amplitude of the 530-GHz mode is small and overshadowed by the 560- and 750-GHz mode to be observed. In a stark contrast to 2.4-THz phonon mode, sub-THz modes exhibit stronger temperature dependences (the frequencies decrease drastically with the increase in temperature), as shown in Fig. 4(b) within the FWHM of the FFT spectra in Fig. 4(a). The comparison of the temperature dependences of the frequencies to those of the previously reported IR transmission spectra of BiFeO₃ [22–26] shows that the 530-, 560-, and 740-GHz modes are consistent with the $\Psi_1^{(2)}$, $\Psi_1^{(1)}$, and $\Phi_2^{(1,2)}$ magnon modes, respectively.

To elucidate the selection rule for the excitation and detection of coherent magnons, we performed time-resolved measurements in X - and Y -probe configurations at various pump azimuth angles θ , as shown in Fig. S5 of SM [44].

Figures 5(a) and 5(b) show the signed amplitudes of the $\Psi_1^{(1)}$ and $\Phi_2^{(1,2)}$ modes in the X - and Y -probe configurations, respectively. Both modes exhibit the same dependence on the pump azimuth [i.e., $\cos(2\theta)$ and $\sin(2\theta)$ in the X - and Y -probe configurations, respectively]. As stated above, the amplitude of the $\Psi_1^{(2)}$ mode is small. However, a FFT analysis at selected pump azimuth reveals the appearance of this mode at $\theta = 0$ and $\pi/2$ in the X -probe configuration and at $\theta = \pm\pi/4$ in the Y -probe configuration as shown in Figs. 5(c) and 5(d). Therefore, the selection rule for the $\Psi_1^{(2)}$ mode is similar to those for the $\Psi_1^{(1)}$ and $\Phi_2^{(1,2)}$ modes.

Further, we use the phenomenological theory of ISRS to correlate the experimental results for BiFeO₃. In principle, the light-matter interaction can be expressed by the interaction free-energy density [46,47],

$$F = -\frac{1}{16\pi} \epsilon_{ij}(\mathbf{M}, \mathbf{L}) \mathcal{E}_i \mathcal{E}_j^*, \quad (1)$$

where $\epsilon_{ij}(\mathbf{M}, \mathbf{L})$ is the dielectric tensor, a complex function of \mathbf{M} and \mathbf{L} (we define a ferromagnetic vector $\mathbf{M} = \mathbf{M}_1 + \mathbf{M}_2$ and antiferromagnetic vector $\mathbf{L} = \mathbf{M}_1 - \mathbf{M}_2$) by using the two sublattice magnetizations \mathbf{M}_1 and \mathbf{M}_2 , and \mathcal{E}_i is the time-dependent amplitude of the electric field $E_i(t) = \text{Re}[\mathcal{E}_i(t) \exp(i\omega t)]$.

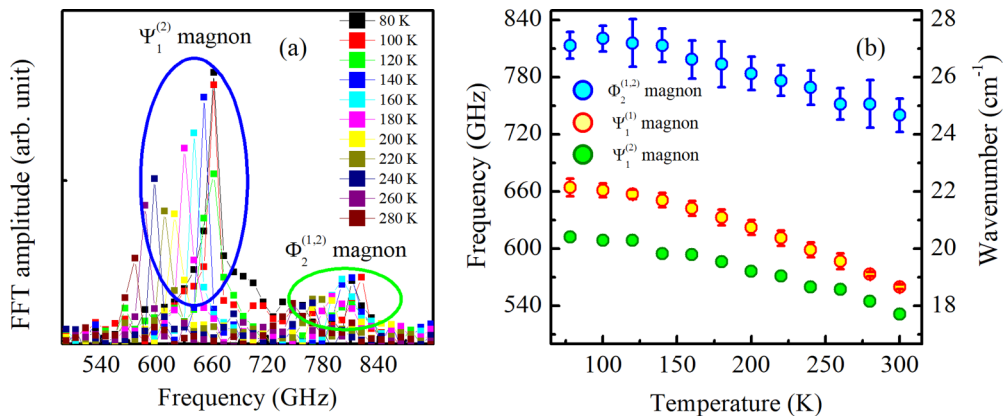


FIG. 4. (a) FFT amplitude spectra of the sub-THz oscillations at various temperatures. (b) Variations of $\Psi_1^{(2)}$ (530 GHz), $\Psi_1^{(1)}$ (560 GHz), and $\Phi_2^{(1,2)}$ (740 GHz) magnon modes as functions of the measurement temperature. Here, the blue and red lines represent the FWHM of the $\Phi_2^{(1,2)}$ and $\Psi_1^{(1)}$ magnon modes, respectively.

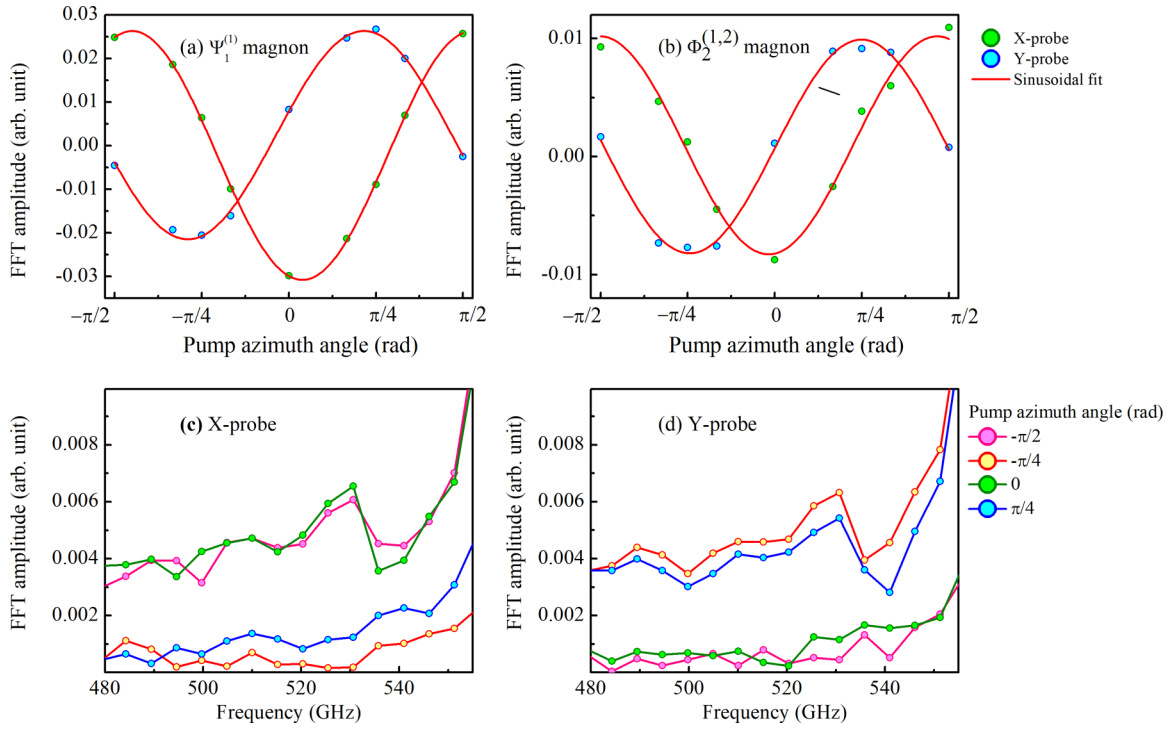


FIG. 5. Dependences of the signed amplitudes of the (a) $\Psi_1^{(1)}$ and (b) $\Phi_2^{(1,2)}$ magnon modes measured in the X-probe (green circles) and Y-probe (blue circles) configurations as functions of the pump azimuth angle ($T = 300$ K). The solid red curves represent the sinusoidal fits of the experimental data. FFT amplitude spectra of 530-GHz $\Psi_1^{(2)}$ magnon mode at selected pump azimuth angles in the (c) X- and (d) Y-probe configurations.

First, we expand the dielectric tensor $\epsilon_{ij}(\mathbf{M}, \mathbf{L})$ into a power series of \mathbf{M} and \mathbf{L} [46,48],

$$\begin{aligned} \epsilon_{ij} = & \epsilon_{ij}^{(0)} + ik_{ijk}^M M_k + ik_{ijk}^L L_k + g_{ijkl}^{MM} M_k M_l \\ & + g_{ijkl}^{LL} L_k L_l + p_{ijkl}^{ML} M_k L_l, \end{aligned} \quad (2)$$

where k_{ijk} and g_{ijkl} are the magneto-optical coupling constants, which mediate the inverse Faraday effect and inverse Cotton-Mouton effect (ICME), respectively [46]. As $|\mathbf{M}| \ll |\mathbf{L}|$, we can neglect the quadratic term of \mathbf{M} , i.e., the fourth term in the right-hand side of Eq. (2). Considering the $3m$ symmetry, the components of k^M , k^L , g^{LL} , and p^{ML} are shown in SM [44] (see also Ref. [49] therein). The antiferromagnetic vector \mathbf{L} exhibits a cycloid rotation in the xz plane [35]. For simplicity, we separately consider the x and z components of \mathbf{L} . We also neglect local ferromagnetic components [50]. We start with $\mathbf{L} = (L_x, 0, 0)$, and thus the components of $\epsilon_{ij}(\mathbf{M}, \mathbf{L})$ can be expressed as in Table I.

TABLE I. Dielectric tensor components of BiFeO₃ for $\mathbf{L} = (L_x, 0, 0)$.

Tensor element	Ψ mode ($m_z, l_y \neq 0$)	Φ mode ($m_y, l_z \neq 0$)
ϵ_{xx}	$p_3 L_x m_z$	$p_4 L_x m_y$
ϵ_{yy}	$-p_3 L_x m_z$	$-p_2 L_x m_y$
ϵ_{xy}	$g_9 L_x l_y + ik_2 m_z$	$g_4 L_x l_z$
ϵ_{yx}	$g_9 L_x l_y - ik_2 m_z$	$g_4 L_x l_z$

In the table, \mathbf{m} and \mathbf{l} are the dynamic components of \mathbf{M} and \mathbf{L} , respectively. The out-of-plane (Ψ) and in-plane (Φ) cyclone modes have m_z , l_y and m_y , l_z nonzero dynamic components, respectively [36,51,52].

We define the effective magnetic fields for \mathbf{m} and \mathbf{l} as

$$\mathbf{H}^{\text{eff}} = -\frac{\partial F}{\partial \mathbf{m}}, \quad (3)$$

$$\mathbf{h}^{\text{eff}} = -\frac{\partial F}{\partial \mathbf{l}} \quad (4)$$

respectively. We define $\mathcal{E}_x(t) = \mathcal{E}_0(t) \cos \theta$ and $\mathcal{E}_y(t) = \mathcal{E}_0(t) \sin \theta \exp(i\varphi)$ ($0 \leq \theta < \pi$, $-\frac{\pi}{2} \leq \varphi < \frac{\pi}{2}$). As the pulse widths of the femtosecond lasers are significantly smaller than the period of the magnon oscillations, $[\mathcal{E}_0(t)]^2 \rightarrow I_0 \delta(t)$, where $I_0 \equiv \int [\mathcal{E}_0(t)]^2 dt$. By using Table I, we can express the components of \mathbf{H}^{eff} and \mathbf{h}^{eff} as

$$H_x^{\text{eff}} = 0, \quad (5)$$

$$H_y^{\text{eff}} = \frac{1}{16\pi} I_0 \delta(t) (p_4 L_x \cos^2 \theta - p_2 L_x \sin^2 \theta), \quad (6)$$

$$H_z^{\text{eff}} = \frac{1}{16\pi} I_0 \delta(t) \{p_3 L_x \cos(2\theta) + k_2 \sin \varphi \sin(2\theta)\}, \quad (7)$$

$$h_x^{\text{eff}} = 0, \quad (8)$$

$$h_y^{\text{eff}} = \frac{1}{16\pi} I_0 \delta(t) g_9 L_x \cos \varphi \sin(2\theta), \quad (9)$$

$$h_z^{\text{eff}} = \frac{1}{16\pi} I_0 \delta(t) g_4 L_x \cos \varphi \sin(2\theta). \quad (10)$$

The equations of motions of \mathbf{m} and \mathbf{l} are [53]

$$\frac{d\mathbf{m}(t)}{dt} = -\frac{1}{2}\gamma(\mathbf{M} \times \mathbf{H}^{\text{eff}} + \mathbf{L} \times \mathbf{h}^{\text{eff}}), \quad (11)$$

$$\frac{d\mathbf{l}(t)}{dt} = -\frac{1}{2}\gamma(\mathbf{M} \times \mathbf{h}^{\text{eff}} + \mathbf{L} \times \mathbf{H}^{\text{eff}}). \quad (12)$$

We neglect the damping terms because they are not important in the analysis. Further, we integrate Eqs. (11) and (12) around $t = 0$, which yields

$$\Delta m = \frac{1}{32\pi}\gamma I_0 L_x^2 [g_4 \cos \varphi \sin(2\theta)\hat{y} - g_9 \cos \varphi \sin(2\theta)\hat{z}], \quad (13)$$

$$\begin{aligned} \Delta l = & \frac{1}{32\pi}\gamma I_0 \left[\{p_3 L_x^2 \cos(2\theta) + k_2 L_x \sin(2\theta) \sin \varphi\} \hat{y} \right. \\ & \left. + \left\{ p_9 L_x^2 \cos(2\theta) - \frac{p_2 - p_4}{2} L_x^2 \right\} \hat{z} \right]. \quad (14) \end{aligned}$$

For a linearly polarized light, $\varphi = 0$, and thus Eqs. (13) and (14) can be simplified to

$$\Delta m^{\text{lin}} = \frac{1}{32\pi}\gamma I_0 L_x^2 [g_4 \sin(2\theta)\hat{y} - g_9 \sin(2\theta)\hat{z}], \quad (15)$$

$$\begin{aligned} \Delta l^{\text{lin}} = & \frac{1}{32\pi}\gamma I_0 \left[p_3 L_x^2 \cos(2\theta)\hat{y} \right. \\ & \left. + \left\{ p_9 L_x^2 \cos(2\theta) - \frac{p_2 - p_4}{2} L_x^2 \right\} \hat{z} \right]. \quad (16) \end{aligned}$$

We define exchange enhancement factors $A_\Psi = \frac{l_y}{m_z}$ and $A_\Phi = \frac{l_z}{m_y}$; usually, A_Ψ and $A_\Phi \gg 1$ [46,47,54]. At $t = 0$, m_y and m_z have finite values, and thus are proportional to $\cos(\omega_\Psi t)$ and $\cos(\omega_\Phi t)$, respectively, whereas l_y and l_z oscillate with a phase difference of $\pi/2$ and exhibit $\sin(\omega_\Psi t)$ and $\sin(\omega_\Phi t)$ dependences, respectively,

$$m_y^{\text{lin}}(t) = \frac{1}{32\pi}\gamma I_0 g_4 L_x^2 \sin(2\theta) \cos(\omega_\Phi t), \quad (17)$$

$$m_z^{\text{lin}}(t) = -\frac{1}{32\pi}\gamma I_0 g_9 L_x^2 \sin(2\theta) \cos(\omega_\Psi t), \quad (18)$$

$$l_y^{\text{lin}}(t) = \frac{1}{32\pi}\gamma I_0 A_\Psi g_9 L_x^2 \sin(2\theta) \sin(\omega_\Psi t), \quad (19)$$

$$l_z^{\text{lin}}(t) = \frac{1}{32\pi}\gamma I_0 A_\Phi g_4 L_x^2 \sin(2\theta) \sin(\omega_\Phi t). \quad (20)$$

As the coupling between the light and magnetization is described solely by the tensor g , we confirm that both modes are originated from the ICME. The changes in probe polarization, i.e., ellipticity ($\Delta\eta$), can be expressed as

$$\begin{aligned} \Delta\eta_X \propto & \epsilon_{xx} - \epsilon_{yy} \propto p_3 L_x m_z + p_9 L_x m_y \\ & \propto p_3 g_9 L_x^3 \sin(2\theta) \cos(\omega_\Psi t) + p_9 g_4 L_x^3 \sin(2\theta) \\ & \times \cos(\omega_\Phi t) \sim (A_{\Psi,\Phi})^0, \quad (21) \end{aligned}$$

$$\begin{aligned} \Delta\eta_Y \propto & \epsilon_{xy} + \epsilon_{yx} \propto g_9 L_x l_y + g_4 L_x l_z \\ & \propto A_\Psi (g_9)^2 L_x^3 \sin(2\theta) \sin(\omega_\Psi t) \\ & + A_\Phi (g_4)^2 L_x^3 \sin(2\theta) \sin(\omega_\Phi t). \quad (22) \end{aligned}$$

TABLE II. Selection rules for the excitations of coherent magnons in the X - and Y -probe configurations.

Probe	Ψ	Φ
X	$\cos(2\theta)$	
Y	$\sin(2\theta)$	$\sin(2\theta)$

Further, we consider the second case for $\mathbf{L} = (0, 0, L_z)$. Similarly (the calculations are shown in SM [44]), under the conditions A_Ψ and $A_\Phi \gg 1$, the components of \mathbf{m} and \mathbf{l} can be expressed as

$$m_x^{\text{lin}}(t) = \frac{1}{32\pi}\gamma I_0 g_4 L_z^2 \cos(2\theta) \cos(\omega_\Psi t), \quad (23)$$

$$m_y^{\text{lin}}(t) = -\frac{1}{32\pi}\gamma I_0 g_4 L_z^2 \sin(2\theta) \cos(\omega_\Phi t), \quad (24)$$

$$l_x^{\text{lin}}(t) = \frac{1}{32\pi}\gamma I_0 A_\Phi g_4 L_z^2 \sin(2\theta) \sin(\omega_\Phi t), \quad (25)$$

$$l_y^{\text{lin}}(t) = \frac{1}{32\pi}\gamma I_0 A_\Psi g_4 L_z^2 \cos(2\theta) \sin(\omega_\Psi t). \quad (26)$$

Likewise, we obtain the changes in probe polarization,

$$\begin{aligned} \Delta\eta_X \propto & g_4 L_z l_y + p_1 L_z m_x \\ & \propto A_\Psi (g_4)^2 L_z^3 \cos(2\theta) \sin(\omega_\Psi t) \\ & + p_1 g_4 L_z^3 \cos(2\theta) \cos(\omega_\Psi t) \\ & \approx A_\Psi (g_4)^2 L_z^3 \cos(2\theta) \sin(\omega_\Psi t), \quad (27) \end{aligned}$$

$$\begin{aligned} \Delta\eta_Y \propto & g_4 L_z l_x - p_1 L_z m_y \\ & \propto A_\Phi (g_4)^2 L_z^3 \sin(2\theta) \sin(\omega_\Phi t) \\ & - p_1 g_4 L_z^3 \sin(2\theta) \cos(\omega_\Phi t) \\ & \approx A_\Phi (g_4)^2 L_z^3 \sin(2\theta) \sin(\omega_\Phi t). \quad (28) \end{aligned}$$

By combining Eqs. (21), (22) and (27), (28) and considering the contributions of L_x and L_z , we obtain the selection rules for the magnon modes, as shown in Table II.

According to Table II, for the Ψ modes, the experimental results well match with the theoretical calculations in both X - and Y -probe configurations. However, the experimental and theoretical results for the Φ mode in the X -probe configuration contradict.

In the above analysis, we assumed that BiFeO₃ has a single magnetic domain with a spin cycloidal rotation along the $\hat{\mathbf{x}} \parallel [10\bar{1}]_{\text{pc}}$ direction. Next, we assume that BiFeO₃ has three uniformly distributed domains where the spin cycloidal rotations are along the $[10\bar{1}]_{\text{pc}}$, $[\bar{1}10]_{\text{pc}}$, and $[0\bar{1}1]_{\text{pc}}$ directions [55]. A detailed description with the contribution of all three domains and probe light polarization changes by L_x and L_z under the condition A_Ψ and $A_\Phi \gg 1$ are presented in SM [44] while the resulting selection rules are shown in Table III.

Thus, the consideration of the contributions of all three domains leads to experimental results well matched with the theoretical calculations.

TABLE III. Selection rules for the excitations of coherent magnons in the X - and Y -probe configurations assuming the contributions of multidomains.

Probe	Ψ	Φ
X	$\cos(2\theta)$	$\cos(2\theta)$
Y	$\sin(2\theta)$	$\sin(2\theta)$

IV. CONCLUSIONS

The dynamics of coherent magnon and phonon modes in the multiferroic BiFeO₃ were investigated by time-resolved ellipticity measurements. Spin and lattice oscillations in BiFeO₃ were excited nonthermally by the linearly polarized light. The subsequent detection by the circularly polarized probe, along with the symmetry analysis by considering the contribution of three uniformly distributed magnetic domains, indicated that the ISRS was responsible for the driving of

the out-of-plane ($\Psi_1^{(2)}, \Psi_1^{(1)}$) and in-plane ($\Phi_2^{(1,2)}$) cyclone magnon modes with frequencies of 530, 560, and 740 GHz at 300 K, respectively. These results pave the way for further studies on electromagnon excitations under the influence of electric and magnetic fields, which could reveal the inherent relationship between the ferroelectric and antiferromagnetic orders in THz optomagnonics.

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